# Fun in codimension 4 

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#### Abstract

I discuss some graded ring constructions of algebraic varieties, mostly motivated by work on algebraic surfaces by Horikawa and his followers. My aim, insofar as possible, is to see the geometric constructions as pullbacks of key varieties. The ideal would be to lift case divisions such as Horikawa's $\mathrm{I}, \mathrm{I}_{a}$ and $\mathrm{I}_{b}$ out of the geometry of surfaces and into general theory of codimension 4 Gorenstein ideals, Tom and Jerry unprojections, key varieties and so on.


## 1 The Horikawa quintics revisited

Horikawa's famous paper [H1] studies canonical surfaces with $p_{g}=4, K^{2}=5$, with the case division
Type I $|K|$ is free and embeds to a quintic $X_{5} \subset \mathbb{P}^{3}$;
Type II $|K|$ has a transversal base point $P$, and, after blowing it up, $\varphi_{K}$ defines a double cover to a quadric $Q \subset \mathbb{P}^{3}$, which may be of rank 4 (Type $\mathrm{II}_{a}$ ) or 3 (Type $\mathrm{II}_{b}$ ).
For details, see [H1], [G], [R2].

### 1.1 The curve and the choice of rendition

In Type II, the curve section $C \in\left|K_{X}\right|$ is a genus 6 hyperelliptic curve with a marked Weierstrass point $P \in C$, polarised by a half-canonical divisor $A=5 P=P+2 g_{2}^{1}$. In coordinates $t_{1}, t_{2}$ on $\mathbb{P}^{1}$, with $P=(0,1)$ and $P_{2}, \ldots, P_{14}$ given by $f_{13}\left(t_{1}, t_{2}\right)=0$, the ring $R(C, A)=R(C, P)^{[5]}$ is generated by

$$
\begin{array}{ll}
\text { in degree } 1 & x_{1}=u t_{1}^{2}, \quad x_{2}=u t_{1} t_{2}, \quad x_{3}=u t_{2}^{2}, \\
\text { in degree 2 } & y_{2}=t_{2}^{5},  \tag{1}\\
\text { in degree } 3 & z_{1}=v t_{1}, \quad z_{2}=v t_{2},
\end{array}
$$

where $u^{2}=t_{1}$ and $v^{2}=f_{13}\left(t_{1}, t_{2}\right)$, and related by

$$
\left.\bigwedge^{2} N=0 \quad \text { where } N=\left(\begin{array}{llll}
x_{1} & x_{2} & x_{3}^{2} & z_{1}  \tag{2}\\
x_{2} & x_{3} & y_{2} & z_{2}
\end{array}\right), \quad \text { and } \quad \begin{array}{r}
z_{1}^{2}=\left[t_{1}^{2} f_{13}\right] \\
z_{1} z_{2}
\end{array}\right)=\left[t_{1} t_{2} f_{13}\right],
$$

The square brackets mean I render forms in $t_{1}, t_{2}$ of degree divisible by 5 (in this case, 15) as weighted forms in $x_{1}, x_{2}, x_{3}, y_{2}$ (in this case, of degree 6). The main point is that different possible renditions give rise to different deformation families of $R(C, A)$ and components of the moduli space of $X$, growing out of the apparently harmless substitution $x_{1} x_{3} \mapsto x_{2}^{2}$. The identity $x_{1} x_{3}=x_{2}^{2}$ in $R(C, P)$ becomes a relation in $R(C, A)$, but, once we deform the ring, will only be a congruence modulo deformation parameters.

Every monomial in $t_{1}^{i} t_{2}^{j}$ of degree 13 has $i \geq 3$ or $j \geq 8$ (with a bit to spare), so I can write $f_{13}$ in the form

$$
\begin{equation*}
f_{13}\left(t_{1}, t_{2}\right)=A_{10} t_{1}^{3}-B_{5} t_{2}^{8} \tag{3}
\end{equation*}
$$

Fix once and for all some rendition $\alpha_{4}, b_{2}$ of $A_{10}, B_{5}$; for example, do

$$
\begin{equation*}
x_{1} x_{3} \mapsto x_{2}^{2}, \quad x_{1} y_{2} \mapsto x_{2} x_{3}^{2}, \quad x_{2} y_{2} \mapsto x_{3}^{3} \tag{4}
\end{equation*}
$$

repeatedly to remove all occurrences of $x_{1} x_{3}, x_{1} y_{2}, x_{2} y_{2}$. Then $\left(t_{1}^{2}, t_{1} t_{2}, t_{2}^{2}\right) f_{13}$ in (2) render as:
(A) $\begin{aligned} & a x_{1}-b x_{3}^{4} \\ & a x_{2}-b x_{3}^{2} y_{2} \quad \text { with } a=\alpha x_{1} \text {; or } \\ & a x_{3}-b y_{2}^{2}\end{aligned}$
(B) $\begin{aligned} & \alpha x_{1}^{2}-b x_{3}^{4} \\ & \alpha x_{1} x_{2}-b x_{3}^{2} y_{2}\end{aligned}$
$\alpha x_{2}^{2}-b y_{2}^{2}$
the only difference being $\alpha x_{1} x_{3} \mapsto \alpha x_{2}^{2}$ in the last line. Case A will correspond to Horikawa's Types $\mathrm{II}_{b}$ and I , whereas Case B will correspond to Types $\mathrm{I}_{b}$ and Type $\mathrm{II}_{a}$.

### 1.2 Case B

There is not too much to say about Case B. The roll $x_{1}^{2} \mapsto x_{1} x_{2} \mapsto x_{2}^{2}$ in (5) is quadratic in the rows, which allows me to replace $N$ in (2) by a general matrix, and the last 3 equations as a general quadratic expression evaluated on its rows. The 9 equations are in rolling factors format:

$$
\bigwedge^{2}\left(\begin{array}{llll}
x_{1} & x_{2}^{\prime} & y_{1} & z_{1}  \tag{6}\\
x_{2} & x_{3} & y_{2} & z_{2}
\end{array}\right)=0 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =\alpha x_{1}^{2}-b y_{1}^{2} \\
z_{1} z_{2} & =\alpha x_{1} x_{2}-b y_{1} y_{2} \\
z_{2}^{2} & =\alpha x_{2}^{2}-b y_{2}^{2}
\end{aligned}
$$

with $x_{2}^{\prime}=x_{2}$ and $y_{1}=x_{3}^{2}$. This format allows the quadric to deform to rank 4 (when $x_{2}^{\prime}$ becomes an independent variable). This construction with general $N$ is an anticanonical divisor in a weighted form of Segre $\mathbb{P}^{1} \times \mathbb{P}^{3}$; it is the main bulk construction of codimension 4 Gorenstein ideals having no known interpretation as a Kustin-Miller unprojection.

### 1.3 Case A

Case A in (5) allows me to roll factors $x_{1} \mapsto x_{2} \mapsto x_{3}$ without putting in terms that are explicitly quadratic in the rows of $N$; this depends on the coincidence $n_{12}=n_{21}=x_{2}$ in $N$, which therefore obstructs deforming the quadric $x_{1} x_{3}-x_{2}^{2}$ to rank $>3$. The variable $x_{1}$ appears linearly in 4 equations multiplying $x_{3}, y_{2}, z_{2}$, $a$, and not in the others, so we can eliminate it, and treat the ring as a Kustin-Miller unprojection from the Pfaffians of

$$
M=\left(\begin{array}{cccc}
0 & x_{2} & x_{3}^{2} & z_{1}  \tag{7}\\
& x_{3} & y_{2} & z_{2} \\
& & z_{2} & -b y_{2} \\
& & & -a
\end{array}\right) \quad \text { of weights } \begin{array}{rrr}
0 & & \\
0 & 2 & 3 \\
1 & 3 \\
3 & 3 \\
3 & 5
\end{array}
$$

with unprojection ideal the codimension 4 c.i. $I=\left(x_{3}, y_{2}, z_{2}, a\right)$. The weight 0 of the entry $m_{12}=0$ is noteworthy. Apart from $m_{13}=x_{2}$ and $m_{15}=z_{1}$, every entry of $M$ is in $I$, so we can treat it in the bigger families of $\mathrm{Tom}_{1}$ or Jerry ${ }_{24}$ unprojections.
$\mathbf{J e r}_{24}$ This case depends on both the coincidences $x_{2}=x_{2}$ and $x_{3} \mid n_{13}$, and does not lead to anything new. For Horikawa surfaces, it only gives deformations inside Type $I_{b}$.

The Jerry format (7) requires $m_{12} \in I$, so $M$ keeps its 0 . It also requires $m_{14}=x_{3}^{2}$ to remain in the ideal $I=\left(x_{3}, y_{2}, z_{2}, A\right)$, thus only allowing $x_{3}^{2}$ to change by adding multiples of $x_{2} x_{3}$ or $y_{2}$; these can be nullified by column operations, so this entry also does not change. The entry $m_{35}=-b y_{2}$ is a free entry, so treat it as a token $B$. After this, the pivot $m_{24}=y_{2}$ can be projected out, and the deformation family calculated as a parallel unprojection. It is a variant on rolling factors:

$$
\bigwedge^{2}\left(\begin{array}{llll}
x_{1} & x_{2} & x_{3}^{2} & z_{1}  \tag{8}\\
x_{2} & x_{3} & y_{2} & z_{2}
\end{array}\right)=0 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =a x_{1}-x_{2} x_{3} B \\
z_{1} z_{2} & =a x_{2}-x_{3}^{2} B \\
z_{2}^{2} & =a x_{3}-y_{2} B
\end{aligned}
$$

Tom $_{1}$ This is where the special hyperelliptic case $\mathrm{I}_{b}$ deforms up to a general quintic hypersurface. The original ring and its general $\mathrm{Tom}_{1}$ deformation are the Pfaffians of the following extrasymmetric $6 \times 6$ matrixes:

$$
M_{0}=\left(\begin{array}{ccccc}
0 & x_{3}^{2} & x_{1} & x_{2} & z_{1}  \tag{9}\\
& y_{2} & x_{2} & x_{3} & z_{2} \\
& & z_{1} & z_{2} & a \\
& & & 0 & b x_{3}^{2} \\
& & & & b y_{2}
\end{array}\right) \quad \mapsto \quad M_{\lambda}=\left(\begin{array}{ccccc}
\lambda & y_{1} & x_{1} & x_{2} & z_{1} \\
& y_{2} & x_{2} & x_{3} & z_{2} \\
& & z_{1} & z_{2} & a \\
& & & \lambda b & b y_{1} \\
& & & & b y_{2}
\end{array}\right)
$$

where $\lambda, y_{1}, b, a$ are indeterminates of weight $0,2,2,5$. Write

$$
\begin{equation*}
\mathcal{C} \mathcal{V} \subset \mathbb{A}_{\left\langle x_{1} \ldots 3, y_{1}, y_{2}, b, z_{1}, z_{2}, a\right\rangle}^{9} \times \mathbb{A}_{\lambda}^{1} \tag{10}
\end{equation*}
$$

for the key variety defined by the $4 \times 4$ Pfaffians of $M_{\lambda}$, the affine cone over the weighted projective variety $\mathcal{V} \subset \mathbb{P}\left(1^{3}, 2^{3}, 3^{2}, 5\right) \times \mathbb{A}^{1}$.

The fibre $\mathcal{C} \mathcal{V}_{\lambda \neq 0}$ is just a copy of $\mathbb{A}_{\left\langle x_{1} \ldots, y_{1}, y_{2}\right\rangle}^{5}$ cunningly set up to degenerate as $\lambda \rightarrow 0$ to the codimension 4 variety $\mathcal{C} \mathcal{V}_{0}$ given by the Pfaffians of $M_{0}$.

Proposition 1.1 Assume first that $\lambda$ is invertible; then $\mathcal{C} \mathcal{V}_{\lambda}$ is the graph over $\mathbb{A}_{\left\langle x_{1} \ldots, y_{1}, y_{2}\right\rangle}^{5}$ of the functions $b, z_{1}, z_{2}$, a defined by four of the Pfaffians $\mathrm{Pf}_{12 . i j}$ :

$$
\begin{align*}
& -\lambda z_{1}=x_{1} y_{2}-x_{2} y_{1}, \quad-\lambda z_{2}=x_{2} y_{2}-x_{3} y_{1} \\
& -\lambda a=y_{2} z_{1}-y_{1} z_{2} \quad \text { and } \quad \lambda^{2} b=x_{1} x_{3}-x_{2}^{2} \tag{11}
\end{align*}
$$

after this, the remaining Pfaffian equations hold as identities. Therefore the fibre $\mathcal{C} \mathcal{V}_{\lambda \neq 0} \cong \mathbb{A}^{5}$ and $\mathcal{V}_{\lambda \neq 0} \cong \mathbb{P}^{4}\left(1^{3}, 2^{2}\right)$.

When $\lambda=0$, the variety $\mathcal{C} \mathcal{V}_{0}$ is given by

$$
\bigwedge^{2}\left(\begin{array}{llll}
x_{1} & x_{2} & y_{1} & z_{1}  \tag{12}\\
x_{2} & x_{3} & y_{2} & z_{2}
\end{array}\right)=0 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =a x_{1}-b y_{1}^{2} \\
z_{1} z_{2} & =a x_{2}-b y_{1} y_{2} \\
z_{2}^{2} & =a x_{3}-b y_{2}^{2}
\end{aligned}
$$

It is an affine Gorenstein codimension 4 variety that can be viewed in an obvious way as an anticanonical divisor in a scroll.

Remark 1.2 The affine cone $\mathcal{C} \mathcal{V}_{0}$ also has a parametric form that displays it as a simple birational transformation away from the quotient of $\mathbb{A}_{\left\langle b, s_{1}, s_{2}, u, v\right\rangle}^{5}$
by the $\mu_{2}$ action $\frac{1}{2}(0,1,1,1,1)$. Indeed, the quantities

$$
\begin{aligned}
x_{1}=s_{1}^{2}, \quad x_{2}=s_{1} s_{2}, \quad x_{3}=s_{2}^{2} \\
y_{1}=s_{1} u, \quad y_{2}=s_{2} u, \\
z_{1}=s_{1} v, \quad z_{2}=s_{2} v,
\end{aligned} \quad \text { and } a=v^{2}+b u^{2}
$$

satisfy the relations (12) identically. In straight projective space $x_{1}, \ldots, z_{2}$ would be coordinates on the scroll $\mathbb{P}_{\mathbb{P}^{1}}(2,1,1)$, the blowup of $\mathbb{P}^{1} \subset \mathbb{P}^{3}$ or the projection of $v_{2}\left(\mathbb{P}^{3}\right)$ from a conic.

### 1.4 Embedded degeneration of quintic curves

The key variety $\mathcal{C V}$ (10) describes degenerations of quintics hypersurfaces, starting from the degeneration of a plane quintic curve to the nonsingular hyperelliptic curve $C_{0} \subset \mathbb{P}(1,1,1,2,3,3)$ as described in Griffin [G]. The simplifying feature here is that it is contained as a complete intersection inside the key variety $\mathcal{V}$.

Consider the complete intersection

$$
\begin{equation*}
\left(b=B_{2}, y_{1}=Y_{2} a=A_{5}\right) \subset \mathcal{V} \subset \mathbb{P}\left(1^{3}, 2^{3}, 3^{2}, 5\right) \times \mathbb{A}^{1} \tag{14}
\end{equation*}
$$

where $B_{2}\left(x_{1 . .3}, y_{2}\right), Y_{2}\left(x_{1 . .3}, y_{2}\right), A_{5}\left(x_{1.3}, y_{2}\right)$ are general forms in $x_{i}, y_{2}$ of the stated weights. For $\lambda \neq 0$, the fibre $\mathcal{V}_{\lambda}$ is $\mathbb{P}(1,1,1,2,2)$, and the two quadratic equations in (14) eliminate $y_{1}, y_{2}$ (because $B_{2}$ contains $y_{2}$ ), so $C_{\lambda}$ is a general plane quintic $C_{5} \subset \mathbb{P}_{\left\langle x_{1 \ldots 3}\right\rangle}^{2}$. For $\lambda=0$, the equations of $\mathcal{V}_{0}$ with the specialisation of $y_{1}, b, a$ are essentially the hyperelliptic equations we started from in (1).

### 1.5 Generalities on regular pullbacks

The above embedded treatment inside $\mathcal{V}$ works for plane quintic curves, but not for higher dimensional quintic hypersurfaces, for example because $h^{0}\left(V_{\lambda}, \mathcal{O}(1)\right)=3$ whereas we need $h^{0}(X, A)=n+2$. The more general notion is regular pullback from a key variety; I explain this briefly for completeness.

By definition, a key variety is an affine variety $W \subset \mathbb{A}^{N}$ that I wish to treat as a key variety (in other words, it is a psychological state); as usual, write $k[W]=k\left[x_{1 \ldots . .}\right] / I_{W}$ for its affine ideal and coordinate ring. $W$ might be, say, the affine cone $a \operatorname{Grass}(2,5) \subset \bigwedge^{2} \mathbb{C}^{5}$ over the Plücker embedding of

Grass $(2,5)$ of [CR], with equations the $4 \times 4$ Pfaffians of a generic $5 \times 5$ skew matrix, or the extrasymmetric variety $\mathcal{C} \mathcal{V} \subset \mathbb{A}^{9} \times \mathbb{A}_{\lambda}^{1}$ of (10), or the origin $0 \in \mathbb{A}^{n}$ defined by the regular sequence $x_{1 \ldots n}$.

Given an ambient ring $R$ (either regular local or polynomial and graded in positive degrees), and a morphism $\varphi: \operatorname{Spec} R \rightarrow \mathbb{A}^{N}$ to the ambient space of $W$, take the pullback or scheme theoretic inverse image $\varphi^{-1} W \subset \operatorname{Spec} R$, and require it to be a regular pullback in the sense of Proposition 1.3. The morphism $\varphi$ specifies values $\varphi^{*}\left(x_{i}\right)=X_{i} \in R$; the pullback is then defined by the ideal $\varphi^{*} I_{W} \subset R$, in other words, by substituting elements $X_{i} \in R$ for $x_{i}$ in the equations of $W$. It is the same thing as the intersection with the graph of $\varphi$

$$
\begin{equation*}
\Gamma_{\varphi} \cap(\operatorname{Spec} R \times W) \subset \operatorname{Spec} R \times \mathbb{A}^{N} \tag{15}
\end{equation*}
$$

the graph $\Gamma_{\varphi}$ is of course the complete intersection cut out by the equations $X_{i}=x_{i}$ for $i=1, \ldots, n$.

## Proposition 1.3 Equivalent conditions:

(i) $X_{i}-x_{i}$ for $i=1, \ldots, n$ form a regular sequence for $\operatorname{Spec} R \times W$.
(ii) The resolution complex of $W$ remains exact on pulling back to Spec $R$. Assume also that $W$ is Cohen-Macaulay; then (i) and (ii) are equivalent to (iii) $\varphi^{-1} W$ has the expected dimension, that is, $\operatorname{codim} \varphi^{-1} W=\operatorname{codim} W$.

In my case (10), I substitute specific values $X_{1}, \ldots, A \in R$ for the variables $x_{1}, \ldots, a$ of $\mathcal{C} \mathcal{V}$ into the extrasymmetric matrix $M_{\lambda}$ of (9), and use the resulting Pfaffians to generate an ideal of $R$.

Even though I am mainly interested in projective varieties and graded rings, the construction itself works on the level of affine cones: $\varphi$ is usually homogeneous (equivariant for appropriate $\mathbb{G}_{m}$ actions), but the induced map $\varphi: \operatorname{Proj} R \rightarrow \mathbb{P}(W)$ need not be a morphism.

### 1.6 Application to Horikawa quintics

We saw that the halfcanonical linear system $A=g_{5}^{2}$ of a nonsingular plane quintic $C_{5}$ can acquire a base point and become $A=P+2 g_{2}^{1}$. The extrasymmetric format (9) also allows the polarising $|\mathcal{O}(1)|$ of a quintic $n$-fold $V_{5}^{n} \subset \mathbb{P}^{n+1}$ to acquire a transverse base point and $\varphi_{\mathcal{O}(1)}$ to degenerate to a double cover of a rank 3 quadric $Q$, while $V_{5}$ remains nonsingular in codimension 2.

Definition 1.4 A numerical quintic is a projective $n$-fold $X$ with at worst terminal singularities, polarised by an ample Cartier divisor $A$, such that

$$
\begin{equation*}
K_{X_{0}}=(3-n) A, \quad A^{n}=5, \quad \text { and } \quad h^{0}(A)=n+2 \tag{16}
\end{equation*}
$$

Theorem 1.5 Let $R=k\left[x_{1 \ldots n+2}, y_{2}, z_{1}, z_{2}\right]$ be the graded polynomial ring with wt $x_{i}, y_{2}, z_{i}=1,2,3$. Let $b=B_{2}, y_{1}=Y_{1,2}, a=A_{5}$ be general forms in $x_{i}, y_{2}$ of the stated weights, and write

$$
\begin{equation*}
\mathcal{X} \subset \operatorname{Proj} R \times \mathbb{A}_{\lambda}^{1}=\mathbb{P}^{n+4}\left(1^{n+2}, 2,3^{2}\right) \times \mathbb{A}_{\lambda}^{1} \tag{17}
\end{equation*}
$$

for the variety defined by the $4 \times 4$ Pfaffians of the extrasymmetric matrix $M_{\lambda}$ of (9). It is a flat family $X_{\lambda}$ of projectively Gorenstein codimension 4 varieties parametrised by $\lambda$, and the fibre $X_{\lambda \neq 0}$ is projectively equivalent to $a$ general quintic in $\mathbb{P}^{n+1}$, lifted to $\mathbb{P}\left(1^{n+2}, 2,3^{2}\right)$ by the forms $b, z_{1}, z_{2}$ of (11) (b contains $y_{2}$ ).

When $\lambda=0$, the Pfaffians take the form (12) with $y_{1}=Q$; the subscheme $X_{0} \subset \mathbb{P}^{n+4}\left(1^{n+2}, 2,3^{2}\right)$ defined by these equations is a numerical quintic with singular locus of dimension $n-3$ (empty if $n \leq 2$ ). The linear system $|A|$ has a single transverse base point $P \in$ NonSing $X_{0}$.

The rational map $\varphi_{|A|}$ blows up $P$ and defines a generically 2-to-1 morphism $\widetilde{\varphi}: \widetilde{X}_{0} \rightarrow Q_{3} \in \mathbb{P}^{n+1}$ where $Q_{3}:\left(x_{1} x_{3}=x_{2}^{2}\right)$.

Geometry of $F \subset \mathbb{P}\left(1^{n+2}, 2\right)$ Consider the involution that acts by -1 on $\lambda, z_{1}, z_{2}$ and fixes $x_{i}, y_{2}$ and $b, y_{1}, a$. Each Pfaffian of $M_{\lambda}$ is $\pm$ invariant (compare (9-12)), and this induces an involution on $\mathcal{X}$ that restricts to a "hyperelliptic" involution on $X_{0}$. The quotient morphism $X_{0} \rightarrow F \subset \mathbb{P}\left(1^{n+2}, 2\right)$ given by the free linear system $|\mathcal{O}(2)|$ is a finite double cover of the codimension 2 determinantal $n$-fold $F$ given by

$$
\bigwedge^{2} N=0 \quad \text { where } \quad N=\left(\begin{array}{lll}
x_{1} & x_{2} & y_{1}  \tag{18}\\
x_{2} & x_{3} & y_{2}
\end{array}\right)
$$

(and $y_{1}=Y_{1,2}\left(x_{i}\right)$ general). This $F$ is singular exactly where $N=0$, together with the quasismooth point $P_{y_{2}} \in F$, an isolated $\frac{1}{2}$ orbifold point.

The reader new to all this should concentrate on the surface case $n=2$, which is familiar from [H1], and relates closely to the relative 2-canonical morphism of a genus 2 fibration at a 2-disconnected fibre as described in $[\mathrm{CP}]$ : the 1-canonical image is then the quadric $Q:\left(x_{1} x_{3}=x_{2}^{2}\right) \subset \mathbb{P}_{\left\langle x_{1 \ldots 4}\right\rangle}^{3}$.

The blown up base point of $\left|K_{S}\right|$ maps to the $x_{3}$-axis $L:\left(x_{1}=x_{2}=0\right)$, and has two marked point $Y_{1,2}=0$ (typically, $x_{3}^{2}-x_{4}^{2}=0$ ) that are the essential singularities of the branch locus in Horikawa's treatment.

Write $Q:\left(x_{1} x_{3}=x_{2}^{2}\right) \subset \mathbb{P}_{\left\langle x_{1} \ldots n+2\right\rangle}^{n+1}$ for the $n$-fold quadric of rank 3, the image of $X_{0}$ under $\varphi_{\mathcal{O}(1)}$. The birational map $\beta: Q \rightarrow F$ is given by quadratic forms on $Q$ allowed poles on the fibre $L=\mathbb{P}^{n-1}:\left(x_{1}=x_{2}\right)$ but required to vanish on $L \cap Y_{1,2}$, giving $y_{2}=x_{2} Y_{1} / x_{1}=x_{3} Y_{1} / x_{2}$ in addition to quadratic forms in $x_{1 \ldots n+2}$. Expressed in birational geometry, $\beta$ first blows up the vertex $x_{1}=x_{2}=x_{3}$ to make the $n$-fold scroll $\mathbb{F}\left(2,0^{n-1}\right)$, then blows up the nonsingular quadric $Y_{1}=0$ in the fibre $L$, and finally contracts $L$ to a $\frac{1}{2}$ orbifold point at $P_{y_{2}}$. The career of the locus $x_{1}=x_{2}=x_{3}=0$ is also interesting: it starts life as the vertex $\mathbb{P}^{n-2}$ of the quadric, is blown up to the negative locus $E=\mathbb{P}^{1} \times \mathbb{P}^{n-2}$ of the scroll. At the fibre $L$ it meets the nonsingular quadric $L \cap Y_{1}$, and after the blowup of $Y_{1}$, is contracted to the $\frac{1}{2}$ orbifold point $P_{y_{2}}$ of the divisor $\mathbb{P}^{n-1}\left(1^{n-1}, 2\right)_{\left\langle x_{4} \ldots n+2, y_{2}\right\rangle} \subset F$, given also by $x_{1}=x_{2}=x_{3}=0$.

The branch locus of the double cover $X_{0} \rightarrow F$ consists of the divisor

$$
\begin{equation*}
D:\left(a x_{1}=b Y_{1}^{2}, \quad a x_{2}=b Y_{1} y_{2}, \quad a x_{3}=b y_{2}^{2}\right) \tag{19}
\end{equation*}
$$

together with the $\frac{1}{2}$ point $P_{y_{2}}$; these are disjoint because $y_{2} \in b=B_{2}\left(x_{i}, y_{2}\right)$. To prove Theorem 1.5, I only need to establish that $D$ is nonsingular outside the singularities of $F$.

Conjecture The Type A family is a generic hypersurface if $\lambda \neq 0$. When $\lambda=0$ it is a birational double cover of a quadric of rank 3 , and is singular at a conic in the vertex (e.g. 2 points if $n=3$ ). The Type $B$ family when $x_{2} \neq x_{2}^{\prime}$ is a double cover of a quadric of rank 4, and is singular at a conic in the vertex (e.g., nonsingular if $n=3$, singular at 2 points if $n=4$ ). It would be interesting to know the relation between the topology of the general Type $A$ and the general Type B, e.g., for the Calabi-Yau case. Can do by computer algebra. Should be easy by Bertini. Hypersurface question on affine pieces.

### 1.7 Comparison with Horikawa's treatment

Horikawa divides surfaces with $p_{g}=4, K^{2}=5$ into three families Type $\mathrm{I}, \mathrm{I}_{b}$ and $\mathrm{II}_{a}$, where I and $\mathrm{II}_{a}$ are the irreducible components of moduli, and $\mathrm{II}_{b}$ is in the closure of both, in codimension 1 in each.

My three cases are

Case A with $\lambda \neq 0$ giving Horikawa's family I,
Case A with $\lambda=0$ giving $\mathrm{I}_{b}$,
Case B with $x_{2}=x_{2}^{\prime}$ also giving $\mathrm{I}_{b}$,
Case B with $x_{2} \neq x_{2}^{\prime}$ giving $\mathrm{II}_{a}$.
When $n=1$ the last 3 cases coincide; they form the hyperelliptic locus, which is a codimension 1 subvariety of family I.

Jul 2011, Feb 2012 I understand this better, but the proof is not written. Regular pullback of my key variety $\mathcal{C V}$ of (10) gives a nonsingular $n$-fold quintic hypersurface $Y_{\lambda} \subset \mathbb{P}^{n+1}$ when $\lambda \neq 0$ degenerating to a codimension $4 n$-fold $Y_{0} \subset \mathbb{P}\left(1^{n+2}, 2,3,3\right)$ when $\lambda=0$; in general $Y_{0}$ has a nonsingular point that is a base point of $|\mathcal{O}(1)|$, and its blowup is a double cover of a quadric of rank 3 ; $Y_{0}$ has ordinary double points over a hyperplane section of the vertex of the quadric, that is, a codimension 3 locus. On the other hand, the format (B) gives rise to $n$-folds that are in general nonsingular in codimension 4.

In the case $n=3$ we have a transition from a nonsingular quintic hypersurface $Y_{5}$ to a nonsingular $Y_{2} \subset \mathbb{P}^{8}\left(1^{5}, 2,3,3\right)$ that is birationally a double cover of a quadric of rank 4. The transition passes through a singular $Y_{2}^{\prime} \subset \mathbb{P}^{8}\left(1^{5}, 2,3,3\right)$ that has two ordinary nodes amd is birationally double cover of a quadric of rank 3 .

These 3 -folds have different Betti numbers $B_{2}$, so are not topologically equivalent.

In the 3 -fold case, I want general case B corresponds to Horikawa fake quintic $Y_{2}$, with a single base point and double cover of quadric of rank 4. Questions: can $Y_{2}$ be nonsingular, and is it diffeo to $Y_{1}=$ quintic in $\mathbb{P}^{4}$. Maybe Pic $Y_{2}=2 \mathbb{Z}$ ?

### 1.8 Other applications

The restriction to $y_{1}=Y_{1,2}$, that is, only one new variable $y_{2}$ in degree 2 , was motivated by the application to quintic hypersurfaces in straight $\mathbb{P}^{n}$. There are many other interesting cases, starting with natural degenerations of hypersurfaces $V_{5} \subset \mathbb{P}\left(1^{n}, 2\right)$ to codimension 4 .

The key variety $\mathcal{C V}$ of (11) has a 3 -parameter family of $\mathbb{C}^{\times}$actions with weights:

$$
\begin{array}{llll}
x_{1} \mapsto n-l & y_{1} \mapsto m & z_{1} \mapsto n+m & b \mapsto 2 n \\
x_{2} \mapsto n & y_{2} \mapsto m+l & z_{2} \mapsto n+m+l & a \mapsto n+2 m+l  \tag{20}\\
x_{3} \mapsto n+l & y_{2} \mapsto n
\end{array}
$$

The determinantal $\bigwedge^{2} N=0$ and its double cover given by (12) apply in other cases. In particular, the same tricks give natural degenerations of K3 and Fano hypersurfaces $V_{5} \subset \mathbb{P}\left(1^{n}, 2\right)$ to codimension 4.

## To finish.

### 1.9 The obstruction

The two deformation families of Case A and Case B are incompatible already at the first infinitesimal level: you can deform to the extrasymmetric format (9) with $\lambda \neq 0$, or you can deform the rank 3 quadric $x_{1} x_{3}-x_{2}^{2}$ to rank 4 , but you can't do both. Compare [R2], Section 5, which calculates Horikawa's obstruction to deformation as $\lambda\left(x_{2}-x_{2}^{\prime}\right)=0$.

## 2 On the BCP construction

Extending Horikawa's work on surfaces with $p_{g}=4, K^{2}=6$, Bauer, Catanese and Pignatelli $[\mathrm{BCP}]$ study deformations of the ring $R\left(C, \frac{3}{2} P\right)$, where $C$ is a hyperelliptic curve of genus 3 and $P \in C$ a Weierstrass point viewed as a $\frac{1}{2}$ orbifold point. Start from the hypersurface

$$
\begin{equation*}
R\left(C, \frac{1}{2} P\right)=k[a, b, c] /\left(c^{2}=f_{7}\left(a^{4}, b\right)\right) \tag{21}
\end{equation*}
$$

with $\operatorname{wt}(a, b, c)=1,4,14$. Its Proj $C_{28} \subset \mathbb{P}(1,4,14)$ has a $\frac{1}{2}$ orbifold point at $P=(1,0,0)$ and ample divisor $A=\frac{1}{2} P$ with $K_{C, \text { orb }}=9 A=2 g_{2}^{1}+\frac{1}{2} P$. The ring $R\left(C, \frac{3}{2} P\right)$ is the third Veronese truncation $R\left(C, \frac{1}{2} P\right)^{[3]}$, and is Gorenstein codimension 4 with generators

$$
x=a^{3}, \quad y=a^{2} b, \quad z=a b^{2}, \quad u=b^{3}, \quad v=a c, \quad w=b c
$$

with $\operatorname{wt}(x, y, z, u, v, w)=1,2,3,4,5,6$ and relations

$$
\bigwedge^{2}\left(\begin{array}{llll}
x & y & z & v  \tag{22}\\
y & z & u & w
\end{array}\right)=0 \quad \begin{array}{ll}
v^{2} & =\left[a^{2} f\right], \\
v w & =[a b f], \\
w^{2} & =\left[b^{2} f\right]
\end{array}
$$

where the square brackets render a form of weight $3 d$ in $a, b$ into a form of weight $d$ in $x, y, z, u$. There are $2 \times 2$ different renditions of (22), that express our ring in four ways as sections of key varieties. The choices are at the two ends of the binary form $f_{7}\left(a^{4}, b\right)$ : terms with high powers of $a$ roll as

$$
\begin{equation*}
a^{3}, a^{2} b, a b^{2} \mapsto x, y, z \quad \text { or as } \quad a^{6}, a^{5} b, a^{4} b^{2} \mapsto x^{2}, x y, y^{2} \tag{23}
\end{equation*}
$$

and at the other end, terms with high powers of $b$ roll as

$$
\begin{equation*}
a^{2} b^{4}, a b^{5}, b^{6} \mapsto z^{2}, z u, u^{2} \quad \text { or as } \quad a^{2} b, a b^{2}, b^{3} \mapsto y, z, u \tag{24}
\end{equation*}
$$

Every monomial $a^{28-4 j} b^{j}$ in $f_{7}\left(a^{4}, b\right)$ has $j \geq 4$ or $28-4 j \geq 4$, so choosing renditions at the two ends of $f$ gives:

$$
\begin{align*}
& v^{2}=A x+C y \quad A x-D z^{2} \quad B x^{2}+C y \quad B x^{2}+D z^{2} \\
& v w=A y+C z \quad \text { or } \quad A y-D z u \quad \text { or } \quad B x y+C z \quad \text { or } \quad B x y+D z u \\
& w^{2}=A z+C u \quad A z-D u^{2} \quad B y^{2}+C u \quad B y^{2}+D u^{2} \tag{25}
\end{align*}
$$

where $A=A_{9}, B=B_{8}, C=C_{8}, D=D_{4}$ are forms of the stated weights in $x, y, z, u$. The four cases in (25) are called (I)-(IV).

Case I This is a double Jerry, see [TJ], Section 8. Two projections eliminate $x$ and $u$ to the codimension 2 c.i. ideal

$$
\begin{equation*}
y w=z v, \quad v w=A y+C z \tag{26}
\end{equation*}
$$

in the product of the ideals $I_{x}=(z, w, A)$ and $I_{u}=(y, v, C)$. The relations (26) deform to the apparently more general form

$$
\left(\begin{array}{lll}
z & w & A
\end{array}\right) M\left(\begin{array}{l}
y  \tag{27}\\
v \\
c
\end{array}\right)=0, \quad\left(\begin{array}{lll}
z & w & A
\end{array}\right) N\left(\begin{array}{l}
y \\
v \\
c
\end{array}\right)=0
$$

with $M$ of weights $\begin{array}{ccc}3 & 0 & -3 \\ 0 & -3 & -6 \\ -3 & -6 & -9\end{array}$ and $N$ of weights $\begin{array}{ccc}6 & & 0 \\ 3 & 0 & 0 \\ 0 & 0 & -3\end{array}$. However, there are not many deformation entries of positive degree in these matrixes, and they can be absorbed by coordinate changes. For example, in $y w-z v+m_{11} z w$, the $m_{11}$ is absorbed by $w \mapsto w+m_{11} z$ or $v \mapsto v-m_{11} y$.

So the variety $W_{8,11} \subset \mathbb{P}(1,2,3,4,5,6,8,9)_{\langle y, z, v, w, C, A\rangle}$ given by (26) and its double unprojection $V \subset \mathbb{P}(1,2,3,4,5,6,8,9)_{\langle x, y, z, u, v, w, C, A\rangle}$ is rigid in these degrees.

Case II Unproject from $x$ to get 5 equations fitting together as the Pfaffians of

$$
\left.\left(\begin{array}{cccc}
0 & y & z & v  \tag{28}\\
& z & u & w \\
& & w & -D u \\
& & & A
\end{array}\right), \quad \text { of weights } \begin{array}{cc}
0 & 2
\end{array}\right)
$$

with unprojection ideal $I_{x}=(z, u, w, A)$. The entry $m_{12}=0$ has weight 0 . Except for $m_{13}, m_{15}$ all the entries are in $I_{x}$, so we can view this as $\mathrm{Tom}_{1}$ or Jerry $_{23}$.

The $\mathrm{Tom}_{1}$ equations give the extrasymmetric format

$$
\left(\begin{array}{ccccc}
0 & z & x & y & v  \tag{29}\\
& u & y & z & w \\
& & v & w & A \\
& & & 0 & D z \\
& & & & D u
\end{array}\right),
$$

We can deform the zero entries $m_{12}$ and $m_{45}$ in (29) to $\lambda$ and $\lambda D$ (with $\lambda$ a variable of degree 0 ), and the entry $m_{24}=z$ to an independent variable $t$ of weight 3 , leading to the matrix

$$
\left(\begin{array}{ccccc}
\lambda & z & x & y & v  \tag{30}\\
& u & y & t & w \\
& & v & w & A \\
& & & \lambda D & D z \\
& & & & D u
\end{array}\right)
$$

When $\lambda \neq 0$ the equations eliminate $v, w, A, D$ to give affine space $\mathbb{A}_{\langle x, y, z, t, u\rangle}^{5}$ or $\mathbb{P}(1,2,3,3,4)_{\langle x, y, z, t, u\rangle}$. Putting back the values of the tokens $A, D$ gives the surface codimension 2 c.i.

$$
\begin{equation*}
S_{4,9}:\left(\lambda^{2} D=x t-y^{2}, \lambda A=z w-u v\right) \subset \mathbb{P}^{4}(1,2,3,3,4)_{\langle x, y, z, t, u\rangle}, \tag{31}
\end{equation*}
$$

where $v=y z-x u$ and $w=z t-y u$. Since $D$ has the same weight as the variable $u$, we can think of $D$ as $\mu u+D^{\prime}$, and for $\mu \neq 0$, this is a general K3 surface $S_{9} \subset \mathbb{P}(1,2,3,3)$ with a built-in degeneration.

Case III Eliminating $u$ gives the ring as the unprojection of the ideal of Pfaffians
in the c.i. ideal $I=(x, y, v, C)$, with the entry $m_{12}=0$ of weight -1 . Except for $m_{24}, m_{25}$, every entry of the matrix is in $I$, so the ring can be viewed either as a $\mathrm{Tom}_{2}$ or Jerry ${ }_{13}$ unprojection.

Because of the -1 , these formats do not allow to lose the $2 \times 3$ minors.
Case IV Every variable appears quadratically in the equations, so there is no naturally occuring unprojection. The deformation family is the matrix format.

## 3 Divisor of odd degree in $v_{2}\left(\mathbb{P}^{2}\right)$

As part of part of the trigonal dichotomy, Castelnuovo, Petri and Mukai tell us that the canonical model $C_{10} \subset \mathbb{P}^{5}$ of a nonhyperelliptic, nontrigonal curve of genus 6 is either a quadric section of a del Pezzo surface $S_{5}$ in the general case, or is the second Veronese embedding $v_{2}\left(C_{5} \subset \mathbb{P}^{2}\right)$ of a plane quintic. In either case there are 6 quadric relations; in the general case these are $5 \times 5$ Pfaffians intersect a quadric hypersurface, leading to a $6 \times 10$ resolution. In the plane quintic case, $C$ needs 3 further cubic equations. The equations of $C$ are $\bigwedge^{2} M=0$ and $\left(A_{1}, A_{2}, A_{3}\right) M=0$ where

$$
M=\left(\begin{array}{lll}
x_{1} & x_{2} & x_{3}  \tag{33}\\
x_{2} & x_{4} & x_{5} \\
x_{3} & x_{5} & x_{6}
\end{array}\right)
$$

corresponds to the Veronese embedding $x_{1}=u^{2}, x_{2}=u v, x_{3}=u w, x_{4}=v^{2}$, $x_{5}=v w, x_{6}=w^{2}$, and the three cubic equations are renditions $\left[u f_{5}\right],\left[v f_{5}\right]$, $\left[w f_{5}\right]$ where $f_{5}(u, v, w)$ defines $C_{5} \subset \mathbb{P}^{2}$.

As before, the key is a seemingly trivial trick with this rendition: observing that every monomial in $u, v, w$ of degree 5 is divisible either by $u$, or by $v^{3}$ or $w^{3}$, I write the equation of $C_{5}$ as $f_{5}=u A+v^{3} B+w^{3} C$ (with $A$ quadratic
in the $x_{i}$ and $B, C$ linear) and the renditions as

$$
\begin{align*}
& u f_{5}=x_{1} A+x_{2} x_{4} B+x_{3} x_{6} C \\
& v f_{5}=x_{2} A+x_{4}^{2} B+x_{5} x_{6} C  \tag{34}\\
& w f_{5}=x_{3} A+x_{4} x_{5} B+x_{6}^{2} C
\end{align*}
$$

Now the set of all 9 equations defining $C$ can be written as $4 \times 4$ Pfaffians of

$$
M=\left(\begin{array}{ccccc}
0 & 0 & x_{1} & x_{2} & x_{3}  \tag{35}\\
& 0 & x_{2} & x_{4} & x_{5} \\
& & x_{3} & x_{5} & x_{6} \\
& & & x_{6} C & -x_{4} B \\
& & & & A
\end{array}\right)
$$

The promising appearance of this as a $6 \times 6$ extrasymmetric matrix is a deception: the top left-hand block cannot become nonzero while preserving the format. I work instead by projecting out $x_{1}$. Then

$$
M=\left(\begin{array}{cccc}
0 & x_{2} & x_{4} & x_{5}  \tag{36}\\
& x_{3} & x_{5} & x_{6} \\
& & x_{6} C & -x_{4} B \\
& & & A
\end{array}\right) \mapsto\left(\begin{array}{cccc}
\lambda & x_{2} & x_{4} & x_{5} \\
& x_{3} & x_{5} & x_{6} \\
& & x_{6} C & -x_{4} B \\
& & & A
\end{array}\right)
$$

is a Jerry ${ }_{45}$ with unprojection ideal $\left(x_{4}, x_{5}, x_{6}, A\right)$. In this format, the three top left entries are free, so I can replace $0 \mapsto \lambda$; since $A$ is a token, I can also project him out to get the equations

$$
\left(\begin{array}{ccc}
x_{3} & -x_{2} & \lambda C  \tag{37}\\
\lambda B & -x_{3} & x_{2}
\end{array}\right)\left(\begin{array}{l}
x_{4} \\
x_{5} \\
x_{6}
\end{array}\right)=0 \quad \Longrightarrow \quad x_{1}\left(\begin{array}{l}
x_{4} \\
x_{5} \\
x_{6}
\end{array}\right)=\left(\begin{array}{c}
x_{2}^{2}-\lambda x_{3} C \\
x_{2} x_{3}-\lambda^{2} B C \\
x_{3}^{2}-\lambda x_{2} B
\end{array}\right)
$$

Thus the equations without $A$ are

$$
\left(\begin{array}{cccc}
x_{1} & x_{2} & x_{3} & \lambda B  \tag{38}\\
& \lambda C & x_{2} & x_{3} \\
& & x_{4} & x_{5} \\
& & & x_{6}
\end{array}\right) .
$$

The final long equation is

$$
\begin{equation*}
x_{1} A+x_{2} x_{4} B+x_{3} x_{6} C-\lambda B C x_{5}=0 . \tag{39}
\end{equation*}
$$

I do not how to write it as a Pfaffian in any useful way. This is characteristic of Jerry.

It is interesting to understand how the deformation $\lambda \neq 0$ allows the 9 equations defining the special $g_{5}^{2}$ curve to pass to just 6 equations defining the general curve. The effect of the $\lambda$ in (38) is to give 5 general equations defining a del Pezzo surface $S_{5}$; the other quadratic equation $f=\lambda A-x_{4} x_{6}+x_{5}^{2}$ defines the general curve $C$ as quadratic section of $S_{5}$. Having done this, the three cubic equations $x_{i} f$ for $i=1,2,3$ become combinations of these 6 .

Wenfei's case: $v_{2}\left(C_{15} \subset \mathbb{P}(1,3,5)\right)$ deforms to $C_{6,6} \subset \mathbb{P}(1,2,3,3)$ The general curve $C_{15} \subset \mathbb{P}(1,3,5)_{\langle u, v, w\rangle}$ is nonsingular, with $K_{C}=6 P$ where $P:\left(v^{5}+w^{3}=0\right) \in \mathbb{P}^{1}(3,5)$. Polarising the same curve by $2 P$ gives the second Veronese $v_{2}\left(C_{15} \subset \mathbb{P}(1,3,5)\right)$; as before, write $x=u^{2}, y=u v, z_{1}=u w$, $z_{2}=v^{2}, s=v w$ and $t=w^{2}$ for the generators, dividing degrees by 2 so that wt $x, y, z_{1}, z_{2}, s, t=1,2,3,3,4,5$; also write $f_{15}=u A_{1}+v A_{2}+w A_{3}$, and render the $A_{i}$ (temporarily) as forms in $x, y, z_{1}, z_{2}, s, t$ of weights $7,6,5$. Then as before, the ring $R(C, 2 P)$ is related by $\bigwedge^{2} M=0$ and $\left(A_{1}, A_{2}, A_{3}\right) M=0$, giving 9 relations of weights $4,5,6,6,7,8,8,9,10$. I can write them as the $4 \times 4$ Pfaffians of

$$
\left.\left(\begin{array}{ccccc}
0 & 0 & x & y & z_{1}  \tag{40}\\
& 0 & y & z_{2} & s \\
& & z_{1} & s & t \\
& & & A_{3} & -A_{2} \\
& & & & A_{1}
\end{array}\right) \quad \text { of weights } \begin{array}{llll} 
& & \\
-1 & 0 & 1 & 2
\end{array}\right)
$$

This extrasymmetric format does not as it stands allow me to deform $C$ to the c.i. $C_{6,6} \subset \mathbb{P}(1,2,3,3)$. As before, a rendition trick is the key: the monomials of $f_{15}(u, v, w)$ not divisible by $v$ are $u^{15}, u^{10} w, u^{5} w^{2}, w^{3}$; therefore, every monomial in $f$ is divisible by $u^{3}$ or $v$ or $w^{3}$, giving the rendition

$$
\begin{equation*}
f_{15}=D u^{3}+B v+E w^{3} \quad \text { with } \quad \text { wt } D, B, E=6,6,0 . \tag{41}
\end{equation*}
$$

The fact that $E$ is a nonzero scalar is the thing that will express $s, t$ as functions of $x, y, z_{1}, z_{2}$ when $\lambda \neq 0$.

Now rewrite the equations not involving $z_{2}$ as the Pfaffians of

$$
\left(\begin{array}{cccc}
0 & x & y & z_{1}  \tag{42}\\
& z_{1} & s & t \\
& & E t & -B \\
& & & D x
\end{array}\right) \mapsto\left(\begin{array}{cccc}
\lambda & x & y & z_{1} \\
& z_{1} & s & t \\
& & E t & -B \\
& & & D x
\end{array}\right) .
$$

This is a Jerry ${ }_{35}$ matrix: the entries in Rows 3 and 5 are in the unprojection ideal $\left(x, z_{1}, t, B\right)$. Almost the same calculation as above give the unprojection equations for $z_{2}$ as the Pfaffians of

$$
\left(\begin{array}{cccc}
\lambda E & x & y & z_{1}  \tag{43}\\
& y & z_{2} & s \\
& & s & t \\
& & & \\
& \lambda D
\end{array}\right) \text { of weights } \begin{array}{rrr}
0 & & \\
0 & 1 & 2 \\
2 & 3 \\
2 & 4 \\
4 & 5 \\
6
\end{array}
$$

together with a long equation for $B z_{2}$. When $\lambda E \neq 0$, these equations eliminate $s, t$

## Full set of equations

$$
\begin{array}{cc}
y^{2}-x z_{2}+\lambda E s & 4 \\
y z_{1}-x s+\lambda E t & 5 \\
x t-z_{1}^{2}+\lambda B & 6  \tag{44}\\
z_{1} z_{2}-y s+\lambda^{2} D E & 6 \\
z_{1} s-y t+\lambda D x & 7
\end{array}
$$

$$
\begin{array}{cc}
D x^{2}+B y+E z_{1} t & 8 \\
s^{2}-z_{2} t+\lambda D y & 8 \\
B z_{2}+D x y+E s t+\lambda D E z_{1} & 9 \\
D x z_{1}+B s+E t^{2} & 10
\end{array}
$$

Notice that if $\lambda=E=1$ (which I can take wlog) then the first 4 equations express $s, t, B, D$ as simple polynomial expressions in $x, y, z_{1}, z_{2}$, and one checks that the remaining 5 equations then hold identically, so that the variety defined by these equation is just the graph over $\mathbb{A}_{\left\langle x, y, z_{1}, z_{2}\right\rangle}^{4}$ of $s, t, B, D$. Substituting general sextics in $x, y, z_{1}, z_{2}$ for $B, D$ defines a complete intersection $C_{6,6} \subset \mathbb{P}(1,2,3,3)$.

Scrap K3 example: take the hypersurface $X_{5} \subset \mathbb{P}(1,1,1,2)$; it is a hypersurface with a $\frac{1}{2}$ orbifold point. If you want to treat it by resolving, you get nonsingular K3 $S$ with fractional divisor $D=B+\frac{1}{2} \Gamma$, where $B$ is ample and $B^{2}=2$, so defines a double cover $S \rightarrow \mathbb{P}^{2}$, and $B \cdot \Gamma=1$, so maps to a bitangent line.

The case $|2 D|$ is in the literature as part of the trigonal dichotomy - the curves in $|2 D|$ have a $g_{5}^{2}$, so the image of $\varphi_{2 D}$ is contained in the Veronese cone $\mathbb{P}(1,1,1,2)=P * v_{2}\left(\mathbb{P}^{2}\right) \subset \mathbb{P}^{6}$. Instead of being an intersection of quadrics, its ideal contains the 6 quadric cones through $v_{2}\left(\mathbb{P}^{2}\right)$ and three cubics corresponding to the rendered products $\left[x_{i} g_{5}\right]$.

The ring $R(X, 2 D)$ is the following codimension 4 structure with $9 \times 16$ resolution. Take the symmetric matrix

$$
M=\left(\begin{array}{lll}
x_{1} & x_{2} & x_{3}  \tag{45}\\
x_{2} & x_{4} & x_{5} \\
x_{3} & x_{5} & x_{6}
\end{array}\right) \text { and set } \bigwedge^{2} M=0 \text { and }\left(\begin{array}{lll}
y_{1} & y_{2} & y_{3}
\end{array}\right) M=0
$$

If you give $x_{i}, y_{i}$ weight one, this is a projectively Gorenstein 4-fold in straight $\mathbb{P}^{8}$ contained in the cone $\mathbb{P}^{2} * v_{2}\left(\mathbb{P}^{2}\right)$. It can be viewed as the degeneration $\lambda \rightarrow 0$ of the extrasymmetric Pfaffian variety

$$
\left(\begin{array}{ccccc}
\lambda y_{3} & -\lambda y_{2} & x_{1} & x_{2} & x_{3}  \tag{46}\\
& \lambda y_{1} & x_{2} & x_{4} & x_{5} \\
& & x_{3} & x_{5} & x_{6} \\
& & & y_{3} & -y_{2} \\
& & & & y_{1}
\end{array}\right) .
$$

Setting $y_{i}$ to be quadratic in the $x_{i}$ gives the canonical curve $C \subset \mathbb{P}^{5}$ or the second Veronese of the $\mathrm{K} 3 X_{5} \subset \mathbb{P}(1,1,1,2)$. The degeneration with $\lambda \rightarrow 0$ in (38) does not explain how to deform it to a general canonical curve or K3 surface. For this, we need to factorise the $y_{i}$ some more.

Let $f_{5}\left(u_{1}, u_{2}, u_{3}\right)$ be the equation of a nonsingular plane quintic. Every monomial in $f$ contains one of $u_{1}^{2}, u_{2}^{2}, u_{3}^{2}$.

So we can render $u_{1} f, u_{2} f, u_{3} f$ as

$$
x_{1} A+x_{2} B+x_{3} C
$$

A similar weighted structure should handle many 2nd Veronese embedding of a hypersurface in $\mathbb{P}(a, b, c)$ or $\mathbb{P}(a, b, c, d)$ with three variables $a, b, c$ of odd weight.

## 4 Horikawa Dicks case

Surfaces with $p_{g}=3, K^{2}=4$. Family $\mathrm{II}_{a}$ : assume the general $C \in\left|K_{S}\right|$ is a hyperelliptic curve of genus 5 polarised by the halfcanonical divisor $A=$ $P_{0}+g_{2}^{1}+P_{\infty}$, where $P_{0}, P_{\infty}$ are Weierstrass points. Take coordinates on $\mathbb{P}^{1}$ so that $P_{0}, P_{\infty} \mapsto 0, \infty$ and $f_{10}\left(t_{1}, t_{2}\right)$ gives the other 10 branch points. Write
$u: \mathcal{O}_{C} \hookrightarrow \mathcal{O}_{C}\left(P_{0}+P_{\infty}\right)$ and $v: \mathcal{O}_{C} \hookrightarrow \mathcal{O}_{C}\left(P_{1}+\cdots+P_{10}\right)$ for the constant sections, so $u^{2}=t_{1} t_{2}$ and $v^{2}=f_{10}\left(t_{1}, t_{2}\right)$. Then $R(C, A)$ is generated by

$$
\begin{align*}
x_{1}=u t_{1}, x_{2}=u t_{2} & \text { in degree } 1, \\
y_{1}=t_{1}^{4}, y_{2}=t_{2}^{4} & \text { in degree } 2,  \tag{47}\\
z_{1}=v t_{1}, \quad z_{2}=v t_{2} & \text { in degree } 3,
\end{align*}
$$

and related by

$$
\operatorname{rank}\left(\begin{array}{llll}
y_{1} & x_{1} & x_{2}^{2} & z_{1}  \tag{48}\\
x_{1}^{2} & x_{2} & y_{2} & z_{2}
\end{array}\right) \leq 1 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =\left[t_{1}^{2} f_{10}\right], \\
z_{1} z_{2} & =\left[t_{1} t_{2} f_{10}\right], \\
z_{2}^{2} & =\left[t_{2}^{2} f_{10}\right],
\end{aligned}
$$

where, as before, the brackets [ ] render the right-hand side as sextics in $x_{i}, y_{i}$. The point is to understand the different ways of doing this.

Remark 4.1 Note that $A=g_{4}^{1}$ on a curve of $g=5$ has Brill-Noether number 1, so imposes 1 condition on the moduli of $C$, and $C, A$ has 11 moduli. The hyperelliptic guy has $2 g-1=9$ moduli, and the trigonal guy with $K_{C}=2\left(g_{3}^{1}+P_{\infty}\right)$ has 10 moduli. The result for curves is that the two fixed points imposes transversal nonsingular divisorial conditions on $C, A$.

### 4.1 Deforming away the base point $P_{0}$

The curve $C$ deforms to lose the fixed point $P_{0}$, so that $A=P_{0}+g_{2}^{1}+P_{\infty} \mapsto$ $g_{3}^{1}+P_{\infty}$. It seems elegant to treat this deformation first in terms of the following bigger variety

$$
\begin{equation*}
V \subset \mathbb{A}_{\left\langle x_{1}, x_{2}, c, y_{1}, y_{2}, D, z_{1}, z_{2}, a, \beta\right\rangle}^{10} \tag{49}
\end{equation*}
$$

defined by

$$
\bigwedge^{2}\left(\begin{array}{rlll}
y_{1} & x_{1} & D & z_{1}  \tag{50}\\
c x_{1} & x_{2} & y_{2} & z_{2}
\end{array}\right)=0 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =A y_{1}+b D^{2} \\
z_{1} z_{2} & =A c x_{1}+b D y_{2}, \\
z_{2}^{2} & =A c x_{2}+b y_{2}^{2}
\end{aligned}
$$

This corresponds to choosing a rendition, and tokenising the features that make possible the subsequent deformation (massage based on hindsight). In detail, any monomial in $f_{10}$ is divisible either by $t_{1}^{2}$ or $t_{2}^{6}$, giving $f_{10}=A t_{1}^{2}+b t_{2}^{6}$, with $A=A_{8}\left(t_{1}, t_{2}\right)$ and $b=b_{4}\left(t_{1}, t_{2}\right)$. I multiply by $t_{1}^{2}, t_{1} t_{2}, t_{2}^{2}$, then substitute

$$
\begin{equation*}
\left(t_{1}^{4}, t_{1}^{3} t_{2}, t_{1}^{2} t_{2}^{2}\right) \mapsto\left(y_{1}, x_{1}^{2}, x_{1} x_{2}\right) \tag{51}
\end{equation*}
$$

in the first summand, and

$$
\begin{equation*}
\left(t_{1}^{2} t_{2}^{6}, t_{1} t_{2}^{7}, t_{2}^{8}\right) \mapsto\left(x_{2}^{4}, x_{2}^{2} y_{2}, y_{2}^{2}\right) \tag{52}
\end{equation*}
$$

in the second. After this, I tokenise $x_{1}^{2}$ as $c x_{1}$ and $x_{2}^{2}=D$.
The resulting set of 9 equations has the following two interpretations as unprojections, where I introduce a deformation parameter $\lambda$ :

$$
y_{1} \cdot\left(x_{2}, y_{2}, z_{2}, A\right) \text { and } \operatorname{Tom}_{1} \text { matrix } \quad M_{1}=\left(\begin{array}{cccc}
\lambda & x_{1} & D & z_{1}  \tag{53}\\
& x_{2} & y_{2} & z_{2} \\
& z_{2} & -b y_{2} \\
& & & A c
\end{array}\right)
$$

and

$$
x_{2} \cdot\left(y_{1}, D, z_{1}, A\right) \text { and Tom } \text { Totrix }_{2} \quad M_{2}=\left(\begin{array}{cccc}
\lambda c & y_{1} & D & z_{1}  \tag{54}\\
& c x_{1} & y_{2} & z_{2} \\
& & z_{1} & -b D \\
& & & A
\end{array}\right) .
$$

The two sets of Pfaffians overlap in two equations for $y_{2} z_{1}$ and $z_{1} z_{2}$; putting them together and coloning out $a$ or $d$ or $y_{2}$ or $z_{1}$ or $z_{2}$ gives the "long equation"

$$
\begin{equation*}
x_{2} y_{1}=c x_{1}^{2}+\lambda^{2} b c . \tag{55}
\end{equation*}
$$

Thus the 9 equations are

$$
\begin{array}{llr}
x_{1} y_{2}=D x_{2}+\lambda z_{2}, & x_{1} z_{2}=x_{2} z_{1}-\lambda b y_{2}, & z_{1}^{2}+A y_{1}+b D^{2}, \\
x_{2} y_{1}=c x_{1}^{2}+\lambda^{2} b c, & y_{1} z_{2}=c x_{1} z_{1}-\lambda b c D, & z_{1} z_{2}+A c x_{1}+b D y_{2},  \tag{56}\\
y_{1} y_{2}=c D x_{1}+\lambda c z_{1}, & y_{2} z_{1}=D z_{2}-\lambda A c, & z_{2}^{2}+A c x_{2}+b y_{2}^{2} .
\end{array}
$$

When $\lambda \neq 0$ these eliminate $z_{2}$, leaving the codimension 3 variety generated by the Pfaffians of

$$
\left(\begin{array}{cccc}
0 & x_{2} & c & y_{2}  \tag{57}\\
& x_{1}^{2}+\lambda^{2} b & y_{1} & \lambda z_{1}+D x_{1} \\
& & D x_{1}-\lambda z_{1} & \lambda^{2} A \\
& & & -D^{2}
\end{array}\right) \quad \text { of weights } \begin{array}{r}
01122 \\
2 \frac{2}{2} 3 \\
34 \\
4
\end{array}
$$

### 4.2 Moving the base point $P_{\infty}$

A parallel interpretation of the original nine equations (48) allows the other base point $P_{\infty}$ to move. I keep $x_{1}^{2}=C$ and $y_{2}^{2}=B$ as tokens (instead of factoring them as $c x_{1}$ and $b y_{2}$ ), but factor the quantities $D$ and $A$ as $D=d x_{2}$ and $A=a y_{1}$. This gives

$$
\bigwedge^{2}\left(\begin{array}{cccc}
y_{1} & x_{1} & d x_{2} & z_{1}  \tag{58}\\
C & x_{2} & y_{2} & z_{2}
\end{array}\right)=0 \quad \text { and } \quad \begin{aligned}
z_{1}^{2} & =a y_{1}^{2}+B d x_{1} \\
z_{1} z_{2} & =a C y_{1}+B d x_{2} \\
z_{2}^{2} & =a C^{2}+B y_{2}
\end{aligned}
$$

The $\mu$ deformation comes from the unprojection interpretations:

$$
y_{2} \cdot\left(x_{1}, y_{1}, z_{1}, B\right) \quad \text { and } \operatorname{Tom}_{2} \text { matrix } \quad M_{1}=\left(\begin{array}{cccc}
\mu & y_{1} & x_{1} & z_{1}  \tag{59}\\
& C & x_{2} & z_{2} \\
& & z_{1} & -B d \\
& & & a y_{1}
\end{array}\right)
$$

and

$$
x_{1} \cdot\left(y_{2}, C, z_{2}, B\right) \quad \text { and Tom }{ }_{1} \text { matrix } \quad M_{2}=\left(\begin{array}{cccc}
\mu d & y_{1} & d x_{2} & z_{1}  \tag{60}\\
& C & y_{2} & z_{2} \\
& & z_{2} & -B \\
& & & a C
\end{array}\right) .
$$

The two sets of Pfaffians overlap in two equations for $y_{1} z_{2}$ and $z_{1} z_{2}$; coloning out gives the "long equation"

$$
\begin{equation*}
x_{1} y_{2}=d x_{2}^{2}+\mu^{2} a d \tag{61}
\end{equation*}
$$

Thus the 9 equations are

$$
\begin{array}{llr}
x_{1} y_{2}=d x_{2}^{2}+\mu^{2} a d, & x_{1} z_{2}=x_{2} z_{1}+\mu a y_{1}, & z_{1}^{2}+a y_{1}^{2}+B d x_{1}, \\
x_{2} y_{1}=C x_{1}+\mu z_{1}, & y_{1} z_{2}=C z_{1}-\mu d B, & z_{1} z_{2}+a y_{1} C+B d x_{2},  \tag{62}\\
y_{1} y_{2}=d C x_{2}+\mu d z_{2}, & y_{2} z_{1}=d x_{2} z_{2}-\mu a d C, & z_{2}^{2}+a C^{2}+B y_{2} .
\end{array}
$$

### 4.3 Putting together the $\lambda$ and $\mu$ deformations

My $\lambda$ and $\mu$ deformation families depend on choices and assumptions that are a priori incompatible if $f_{10}$ has a nonzero term in $t_{1}^{5} t_{2}^{5}$. Ignoring this for the moment, assume that $f_{10}=a_{4} t_{1}^{6}+b t_{2}^{6}$. With a little trial and error, one
checks that the $\lambda$ and $\mu$ deformations (56) and (62) fit together, somewhat miraculously, with only a single $\lambda \mu$ term in the $z_{1} z_{2}$ equation:

$$
\begin{align*}
& x_{1} y_{2}=d x_{2}^{2}+\lambda z_{2}+\mu^{2} a d, \\
& x_{2} y_{1}=c x_{1}^{2}+\lambda^{2} b c+\mu z_{1}, \\
& y_{1} y_{2}=c d x_{1} x_{2}+\lambda c z_{1}+\mu d z_{2}, \\
& x_{1} z_{2}=x_{2} z_{1}-\lambda b y_{2}+\mu a y_{1}, \\
& y_{1} z_{2}=c x_{1} z_{1}-\lambda b c d x_{2}-\mu b d y_{2},  \tag{63}\\
& y_{2} z_{1}=d x_{2} z_{2}-\lambda a c y_{1}-\mu a c d x_{1}, \\
& z_{1}^{2}+a y_{1}^{2}+b d x_{1} y_{2}-\lambda b d z_{2}, \\
& z_{1} z_{2}+a c x_{1} y_{1}+b d x_{2} y_{2}-\lambda \mu a b c d, \\
& z_{2}^{2}+a c x_{2} y_{1}+b y_{2}^{2}-\mu a c z_{1} .
\end{align*}
$$

I assert that setting $\mu$ or $\lambda$ to zero gives back the known $\lambda$ and $\mu$ deformation families, and that these equations define a flat deformation over $\mathbb{A}_{\langle\lambda, \mu\rangle}^{2}$. To check flatness, it is enough to check that the 16 syzygies (66) hold (with $e=0$ ).

Finally, I deal with the missing term in $a_{5} t_{1}^{5} t_{2}^{5}$ in $f_{10}\left(t_{1}, t_{2}\right)$ by setting $f_{10}=a_{4} t_{1}^{6}+e t_{1} t_{2}+b_{4} t_{2}^{6}$ where $e=a_{5} y_{1} y_{2}$, and render it as $L_{7} \mapsto L_{7}+e x_{1}^{2}$, $L_{8} \mapsto L_{8}+e x_{1} x_{2}, L_{9} \mapsto L_{9}+e x_{2}^{2}$ or

$$
\begin{align*}
t_{1}^{2} f_{10} & =a y_{1}^{2}+e x_{1}^{2}+b d x_{1} y_{2}, \\
t_{1} t_{2} f_{10} & =a c x_{1} y_{1}+e x_{1} x_{2}+b d x_{2} y_{2},  \tag{64}\\
t_{2}^{2} f_{10} & =a c x_{2} y_{1}+e x_{2}^{2}+b y_{2}^{2}
\end{align*}
$$

The equations become

$$
\begin{align*}
& L_{1}: x_{1} y_{2}=d\left(x_{2}^{2}+\mu^{2} a\right)+\lambda z_{2}, \\
& L_{2}: x_{2} y_{1}=c\left(x_{1}^{2}+\lambda^{2} b\right)+\mu z_{1} \text {, } \\
& L_{3}: y_{1} y_{2}=c d x_{1} x_{2}+\lambda c z_{1}+\mu d z_{2}-\lambda \mu e \\
& \equiv c\left(d x_{1} x_{2}+\lambda z_{1}\right)+\mu\left(d z_{2}-\lambda e\right) \\
& \equiv d\left(c x_{1} x_{2}+\mu z_{2}\right)+\lambda\left(c z_{1}-\mu e\right), \\
& L_{4}: x_{1} z_{2}=x_{2} z_{1}-\lambda b y_{2}+\mu a y_{1} \text {, }  \tag{65}\\
& L_{5}: y_{1} z_{2}=\left(c z_{1}-\mu e\right) x_{1}-b d\left(\lambda c x_{2}+\mu y_{2}\right), \\
& L_{6}: y_{2} z_{1}=\left(d z_{2}-\lambda e\right) x_{2}-a c\left(\lambda y_{1}+\mu d x_{1}\right) \text {, } \\
& L_{7}: \quad z_{1}^{2}+a y_{1}^{2}+e x_{1}^{2}+b d x_{1} y_{2}-\lambda b\left(d z_{2}-\lambda e\right)=0, \\
& L_{8}: z_{1} z_{2}+a c x_{1} y_{1}+e x_{1} x_{2}+b d x_{2} y_{2}-\lambda \mu a b c d=0, \\
& L_{9}: \quad z_{2}^{2}+a c x_{2} y_{1}+e x_{2}^{2}+b y_{2}^{2}-\mu a\left(c z_{1}-\mu e\right)=0 .
\end{align*}
$$

This set of equations comes neatly from $I_{0}=\left(L_{1}, L_{2}, L_{4}, L_{8}\right)$ (unchanged from (63) except for the unsurprising term $e x_{1} x_{2}$ in $L_{8}$ ) by coloning out
$x_{1} x_{2} y_{1} y_{2}$; its syzygy matrix $M$ is

| $y_{1}$ | $d x_{2}$ | $-x_{1}$ | $-\mu d$ | $\lambda$ | $\cdot$ | $\cdot$ | $\cdot$ | $\cdot$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $c x_{1}$ | $y_{2}$ | $-x_{2}$ | $\lambda c$ | $\cdot$ | $\mu$ | $\cdot$ | $\cdot$ | $\cdot$ |
| $\cdot$ | $-z_{1}$ | $\lambda b$ | $-y_{1}$ | $x_{1}$ | $\cdot$ | $\mu$ | $\cdot$ | $\cdot$ |
| $\lambda b c$ | $-z_{2}$ | $\cdot$ | $-c x_{1}$ | $x_{2}$ | $\cdot$ | $\cdot$ | $\mu$ | $\cdot$ |
| $-z_{1}$ | $\mu a d$ | $\cdot$ | $d x_{2}$ | $\cdot$ | $x_{1}$ | $\cdot$ | $\lambda$ | $\cdot$ |
| $-z_{2}$ | $\cdot$ | $\mu a$ | $y_{2}$ | $\cdot$ | $x_{2}$ | $\cdot$ | $\cdot$ | $\lambda$ |
| $\cdot$ | $-a y_{1}$ | $\cdot$ | $z_{1}$ | $\cdot$ | $-\lambda b$ | $-x_{2}$ | $x_{1}$ | $\cdot$ |
| $b y_{2}$ | $\cdot$ | $\cdot$ | $z_{2}$ | $\mu a$ | $\cdot$ | $\cdot$ | $-x_{2}$ | $x_{1}$ |

$$
\begin{array}{ccccccccc}
. & \mu a e & \cdot & a c y_{1}+e x_{2} & \cdot & -b y_{2} & -\mu a c & z_{2} & -z_{1}  \tag{66}\\
\lambda b e & \cdot & \cdot & -b d y_{2}-e x_{1} & -a y_{1} & \cdot & -z_{2} & z_{1} & -\lambda b d \\
a c y_{1}+e x_{2} & \cdot & \cdot & -\mu a c d & \cdot & z_{2} & \cdot & y_{2} & -d x_{2} \\
-e x_{1} & \cdot & -a y_{1} & d z_{2}-\lambda e & \cdot & -z_{1} & -y_{2} & \cdot & d x_{1} \\
\cdot & -e x_{2} & -b y_{2} & -c z_{1}+\mu e & -z_{2} & \cdot & c x_{2} & \cdot & -y_{1} \\
\cdot & b d y_{2}+e x_{1} & \cdot & \lambda b c d & z_{1} & \cdot & -c x_{1} & y_{1} & \cdot \\
-c z_{1}+\mu e & \cdot & z_{2} & c d x_{2} & -y_{2} & \cdot & \cdot & \cdot & -\mu d \\
\cdot & d z_{2}-\lambda e & -z_{1} & c d x_{1} & \cdot & y_{1} & \lambda c & \cdot & \cdot
\end{array}
$$

(Or ad lib, apply opposite row operation to Rows 1-8 and Rows 9-16, or swap Rows $i$ and $i+8$.) One checks that it satisfies ${ }^{t} M J M=0$ where $J=\left(\begin{array}{cc}0 & I \\ I & 0\end{array}\right)$ is the standard quadratic form, so that essentially the same matrix $M$ also provides the second syzygies, giving the projective resolution

$$
\begin{equation*}
\mathcal{O}_{X} \leftarrow \mathcal{O} \leftarrow P_{1} \leftarrow P_{2} \leftarrow P_{3} \leftarrow P_{4} \leftarrow 0, \tag{67}
\end{equation*}
$$

with

$$
\begin{align*}
& P_{1}=2 \mathcal{O}(-3) \oplus 2 \mathcal{O}(-4) \oplus 2 \mathcal{O}(-5) \oplus 3 \mathcal{O}(-6), \\
& P_{2}=2 \mathcal{O}(-5) \oplus 4 \mathcal{O}(-6) \oplus 4 \mathcal{O}(-7) \oplus 4 \mathcal{O}(-8) \oplus 2 \mathcal{O}(-10),  \tag{68}\\
& P_{4}=\mathcal{O}(-14) \quad \text { and } \quad P_{3}=\mathcal{H o m}\left(P_{2}, P_{4}\right)=P_{2}^{\vee} \otimes \mathcal{O}(-14) .
\end{align*}
$$

### 4.4 Set $\lambda \neq 0$ and eliminate $z_{2}$

If $\lambda$ is invertible, $L_{1}$ gives $z_{2}=\left(\left(x_{2}^{2}+\mu^{2} a\right) d-x_{1} y_{2}\right) / \lambda$, and (65) boil down to the Pfaffians of

$$
\left(\begin{array}{cccc}
\mu & x_{2} & c & y_{2}  \tag{69}\\
& x_{1}^{2}+\lambda^{2} b & y_{1} & \lambda z_{1}+d x_{1} x_{2} \\
& & z_{1} & -a\left(\lambda y_{1}+\mu d x_{1}\right) \\
& & & \lambda e-d z_{2}
\end{array}\right)
$$

the first 4 of which give

$$
\begin{align*}
& c\left(x_{1}^{2}+\lambda^{2} b\right)-x_{2} y_{1}+\mu z_{1}, \\
& \mu a\left(\lambda y_{1}+\mu d x_{1}\right)+x_{2}\left(\lambda z_{1}+d x_{1} x_{2}\right)-\left(x_{1}^{2}+\lambda^{2} b\right) y_{2},  \tag{70}\\
& \mu \lambda e-\mu d z_{2}-\lambda c z_{1}-c d x_{1} x_{2}+y_{1} y_{2}, \\
& \lambda e x_{2}-d x_{2} z_{2}+a c\left(\lambda y_{1}+\mu d x_{1}\right)+y_{2} z_{1},
\end{align*}
$$

whereas

$$
\begin{equation*}
\mathrm{Pf}_{23.45}=\mu a d x_{1} y_{1}+d x_{1} x_{2} z_{1}-\lambda^{2} b d z_{2}-d x_{1}^{2} z_{2}+\lambda a y_{1}^{2}+\lambda z_{1}^{2}+\lambda e\left(x_{1}^{2}+\lambda^{2} b\right) . \tag{71}
\end{equation*}
$$

After subtracting $d x_{1} L_{4}$, this is divisible by $\lambda$ and gives
$L_{7}=\mu^{2} a b d^{2}+b d^{2} x_{2}^{2}+a y_{1}^{2}+z_{1}^{2}+\lambda^{2} b e+e x_{1}^{2}$.
Similarly if $\mu$ is invertible, set $z_{1}=\left(\left(x_{1}^{2}+\lambda^{2} b\right) c-x_{2} y_{1}\right) / \mu$

$$
\left(\begin{array}{cccc}
\lambda & x_{1} & d & y_{1}  \tag{72}\\
& x_{2}^{2}+\mu^{2} a & y_{2} & \mu z_{2}+c x_{1} x_{2} \\
& & z_{2} & -b\left(\mu y_{2}+\lambda c x_{2}\right) \\
& & & \mu e-c z_{1}
\end{array}\right)
$$

The two matrixes have a common Pfaffian 12.45, and (after cancelling $\lambda$ and $\mu$ judiciously), their Pfaffians together generate the ideal (63). Check that

$$
\begin{align*}
-\lambda L[4] & =x_{1} \operatorname{Pf}_{12.34}\left(M_{\lambda}\right)+\operatorname{Pf}_{12.35}\left(M_{\mu}\right),  \tag{73}\\
\mu L[4] & =x_{2} \operatorname{Pf}_{12.34}\left(M_{\mu}\right)+\operatorname{Pf}_{12.35}\left(M_{\lambda}\right),  \tag{74}\\
\lambda \mu L[7] & =\lambda z_{1} \operatorname{Pf}_{12.34}\left(M_{\mu}\right)+\left(\lambda^{2} b+x_{1}^{2}\right) \operatorname{Pf}_{12.45}\left(M_{\lambda}\right)-\left(y_{1}+d x_{1}\right) \operatorname{Pf}_{12.35}\left(M_{\lambda}\right) . \tag{75}
\end{align*}
$$

TO BE CONTINUED

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