

Solutions to exercise sheet 2

$$\boxed{1} \quad \text{For } d=2, \text{ we have } \partial_{x_i} F(\vec{x}) = -\frac{1}{2\pi} \partial_x (\ln(\sqrt{x^2+y^2})) =$$

$$= -\frac{1}{2\pi} \frac{1}{\sqrt{x^2+y^2}} \partial_x (\sqrt{x^2+y^2}) = -\frac{1}{2\pi} \frac{x}{x^2+y^2}.$$

$$\text{So } \partial_x^2 F(\vec{x}) = -\frac{1}{2\pi} \left(\frac{1}{x^2+y^2} - \frac{2x^2}{(x^2+y^2)^2} \right)$$

In the same way, we see

$$\partial_y^2 F(x,y) = -\frac{1}{2\pi} \left(\frac{1}{x^2+y^2} - \frac{2y^2}{(x^2+y^2)^2} \right),$$

and thus

$$\Delta F(x,y) = -\frac{1}{2\pi} \left(\frac{2}{x^2+y^2} - \frac{2(x^2+y^2)}{(x^2+y^2)^2} \right) = 0.$$

For $d \geq 3$, we find

$$\partial_{x_i} F(\vec{x}) = \frac{1}{n(n-2)d(n)} \partial_{x_i} \frac{1}{\left(\sum_{j=1}^n x_j^2\right)^{(n-2)/2}} = \frac{1}{n(n-2)d(n)} \frac{-\left(\frac{n-2}{2}\right) \cdot 2x_j}{\left(\sum_{j=1}^n x_j^2\right)^{\frac{n-2}{2}+1}}$$

$$= \frac{-1}{n d(n)} \frac{x_j}{\left(\sum_{j=1}^n x_j^2\right)^{n/2}}, \quad \text{and}$$

$$\partial_{x_i}^2 F(\vec{x}) = \frac{-1}{n d(n)} \left(\frac{1}{\left(\sum_{j=1}^n x_j^2\right)^{n/2}} - \frac{\frac{n}{2} \cdot 2x_j^2}{\left(\sum_{j=1}^n x_j^2\right)^{\frac{n+2}{2}}} \right)$$

$$\text{So } \Delta F(\vec{x}) = -\frac{1}{n d(n)} \left(\frac{n}{|\vec{x}|^n} - \frac{n |\vec{x}|^2}{|\vec{x}|^{n+2}} \right) = 0.$$

2 The reasoning is again, like last week's:

$$u(\vec{x}) = \mathbb{E}_{\vec{y}_0 = \vec{x}} (u(\vec{y}_0)), \text{ and}$$

$$u(\vec{y}_{J(\vec{x})}) = \mathbb{E}_{\vec{y}_0 = \vec{y}_{J(\vec{x})}} (\dots) = \bar{\Phi}(\vec{y}_{J(\vec{x})}).$$

$$\begin{aligned} \text{So, } 0 &= \mathbb{E}_{\vec{y}_0 = \vec{x}} \left(\int_0^{J(\vec{x})} \mathcal{F}(\vec{y}_s) ds + \underbrace{\bar{\Phi}(\vec{y}_{J(\vec{x})}) - u(\vec{y}_0)}_{= u(\vec{y}_{J(\vec{x})})} \right) \\ &= \mathbb{E}_{\vec{y}_0 = \vec{x}} \left(\int_0^{J(\vec{x})} \mathcal{F}(\vec{y}_s) ds + \int_0^{J(\vec{x})} du(\vec{y}_s) \right). \end{aligned}$$

We know that $du(\vec{y}_s) = \underbrace{\frac{\partial}{\partial t} u(\vec{y}_s)}_{\equiv 0} ds + \mathcal{L}u(\vec{y}_s) ds + \sum_{i,j=1}^n g_{ij}(\vec{y}_s) (\partial_{x_i} u)(\vec{y}_s) dW_s^{(i)},$

with $\mathcal{L}u = \sum_{i=1}^n f_i \partial_{x_i} u + \frac{1}{2} \sum_{i,j,k=1}^n g_{ik} g_{jk} \partial_{x_i} \partial_{x_j} u,$

by the Itô formula. The expectations of the stochastic integrals vanish, so u must solve

$$\begin{aligned} \mathcal{L}u + \mathcal{F} &= 0 \quad \text{on } \mathcal{D} \\ u &= \bar{\Phi} \quad \text{on } \partial\mathcal{D}. \end{aligned}$$

3 Assume that u solves

$$x \cdot \nabla u = \alpha u$$

Define $v: \mathbb{R} \rightarrow \mathbb{R}$ by $v(\lambda) = u(\lambda x)$.

$$\begin{aligned} \text{Then } \frac{d}{d\lambda} v(\lambda) &\stackrel{\text{chain rule}}{=} x \cdot \nabla u(\lambda x) = \frac{1}{\lambda} (\lambda x) \cdot \nabla u(\lambda x) = \\ &\stackrel{\text{PDE}}{=} \frac{1}{\lambda} \alpha u(\lambda x) = \frac{\alpha}{\lambda} v(\lambda). \end{aligned}$$

So v solves the ODE $v' = \frac{\alpha}{\lambda} v$, with initial value

$v(1) = u(1 \cdot x) = u(x)$. The unique solution to this

ODE is $v(\lambda) = v(1) \cdot \lambda^\alpha$, and thus we must have

$$u(\lambda x) = v(\lambda) = v(1) \lambda^\alpha = \lambda^\alpha u(x).$$

4 a) The PDE can be written as

$$b \cdot \nabla u = u \quad \text{with } b = (1, 1).$$

Finding a curve such that $\dot{y}(s) = b$ is easy:

$$y(s) = (s + \tau, s) \quad \text{for any } \tau \in \mathbb{R}.$$

Along y , we find

$$\frac{d}{ds} v(s) = \frac{d}{ds} u(y(s)) = b \cdot \nabla u(y(s)) = u(y(s)) = v(s),$$

so $v(s) = e^s \cdot \text{const.}$ By using

$$v(0) = u(y(0)) = u(r, 0) \stackrel{\text{Boundary cond.}}{=} f(r),$$

we obtain $v(s) = e^s f(r)$.

In the original coordinates, [since $x = s+r$, $y = s \Rightarrow$

$$\Rightarrow s = y \text{ and } r = x - s = x - y], \text{ we have finally}$$

$$u(x, y) = e^y f(x - y).$$

b) The initial data $u(x, x) = 1$ is on one of the curves

$\gamma(s) = (s+r, s)$ from a), namely for $r=0$. We have

$$\frac{d}{ds} v(s) = v(s) \quad \text{and} \quad v(s) = u(\gamma(s)) = u(s, s) = 1 \text{ for all } s.$$

This is impossible (as $v'(s) = v(s)$ only admits $v(s) = c e^s \neq 1$),
so there is no solution.

c) Same as in b), but now we get

$$\frac{d}{ds} v(s) = v(s) \quad \text{and} \quad v(s) = 0. \text{ This has the solution}$$

$v(s) = 0$ for all s . But for all other lines

$\gamma(s) = (s+r, s)$ ($r \neq 0$) we have no prescribed value,

and thus can choose any solution

$$v(s) = c(r) e^s, \text{ where } c(r) \text{ is arbitrary.}$$

So the solution is not unique.